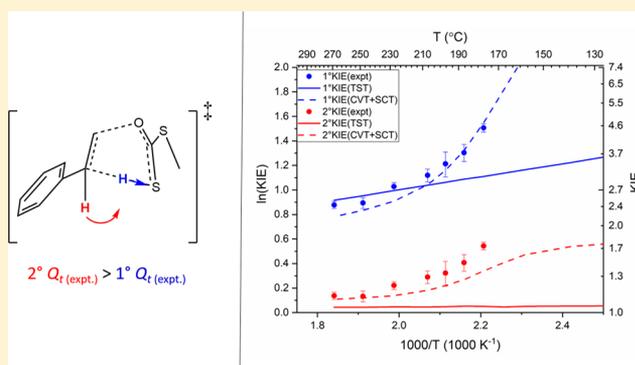


## Quantum Suppression of Intramolecular Deuterium Kinetic Isotope Effects in a Pericyclic Hydrogen Transfer Reaction

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## Supporting Information

**ABSTRACT:** It is generally accepted that hydrogen tunneling enhances both primary and secondary H/D kinetic isotope effects (KIEs) over what would be expected under the assumptions of classical barrier transition. Previous studies have exclusively shown that the effects of tunneling upon primary H/D KIEs have been much larger than those observed for secondary H/D KIEs. Here we present a series of experimental H/D KIE results associated with the Chugaev elimination of methyl xanthate derived from  $\beta$ -phenylethanol over the temperature range of 180 to 290 °C. Intramolecular H/D KIEs computed according to classical transition state theory (TST) are markedly overestimated relative to experimentally measured values. Experimental intermolecular H/D KIEs and direct dynamic calculations based on canonical variational transition state theory (CVT) with small-curvature tunneling correction (SCT) reveal that this result is largely the consequence of extraordinary tunneling enhancement of the *secondary* H/D KIE. This unexpected behavior is examined in the context of other similar hydrogen transfer reactions.



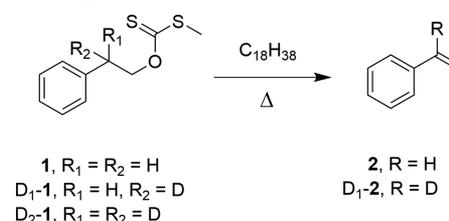
## 1. INTRODUCTION

Hydrogen tunneling has generated a lot of debate within the chemical and biochemical communities.<sup>1–5</sup> This is because hydrogen transfer is important to numerous synthetically useful transformations and the vast majority of biological processes, while the effects of tunneling upon hydrogen transfer remain poorly understood in many reactions. Quantum mechanical tunneling exhibits phenomena that cannot be accounted for using classical transition state theory (TST), which is the most generally applicable theory for chemical kinetics within the classical regime.<sup>6,7</sup> Cogent models exist for proton-coupled electron transfer reactions that yield extremely large primary deuterium KIEs<sup>8,9</sup> and for atom–diatomic collinear reactions.<sup>10,11</sup> Surprisingly, the quantum mechanical behavior of hydrogen remains poorly understood for many heterolytic polyatomic hydrogen transfer reactions. This knowledge gap hinders our ability to understand reactivity and selectivity in reactions significant to both metabolism and chemical synthesis. Improved models for reactions that exhibit significant influence from quantum mechanical tunneling promise to facilitate improved enzyme inhibitors for small molecule drug design and better catalysts for reactions involving hydrogen transfer.<sup>12–14</sup> Before better models can be constructed for reactions that involve substantial quantum behavior, an expanded set of well-interpreted experimental results needs to be presented. This paper is intended to present

a model system in which common metrics for tunneling indicate classical behavior that masks surprising nonclassical behavior.

This work studies the tunneling effect in the intramolecular *syn*- $\beta$ -elimination reaction in *S*-methyl *O*-(2-phenyl)ethyl xanthate **1** (Scheme 1) by examining the temperature dependence of the KIEs through both experiments and calculations. Kinetic isotope effects are often the observable chosen to demonstrate the importance of tunneling in a given reaction.<sup>15–22</sup> Efforts over the past couple of decades have attempted to unite a conceptual view of hydrogen tunneling with both magnitudes and temperature dependencies of KIEs.

## Scheme 1. Chugaev Elimination



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In spite of extensive work, a consensus has not been reached concerning the strange temperature dependencies of primary deuterium KIEs observed in some enzymatic systems.<sup>23–25</sup> Although progress has been made toward a conventional wisdom regarding the temperature dependence of KIEs in tunneling systems, new findings continue to challenge our working models.

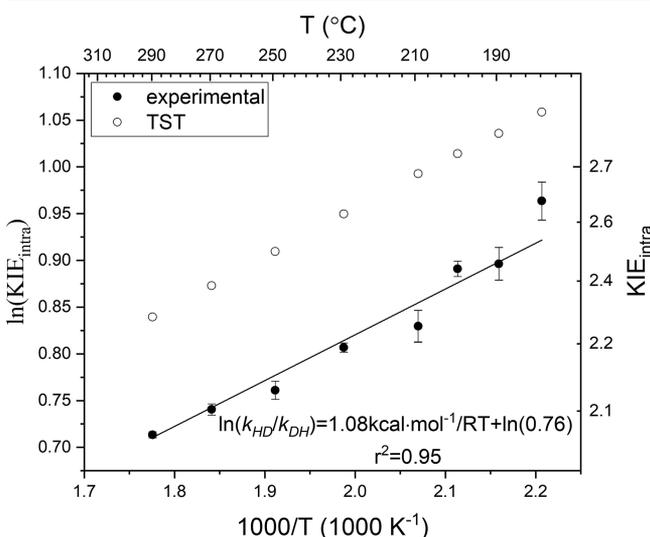
Our experimental results show that the Chugaev elimination of **1** exhibits a deceptively classical Arrhenius plot of the intramolecular H/D KIE that masks intrinsically nonclassical behavior of the primary and secondary KIEs. These results challenge commonly accepted expectations by demonstrating that (1) tunneling can have profound influence upon isotope effects measured at temperatures well in excess of room temperature and (2) that secondary KIEs can be inflated by tunneling to a *greater* degree than primary KIEs. Computational modeling of the Chugaev elimination using canonical variational transition state theory (CVT) in conjunction with small-curvature tunneling (SCT) correction for quantum barrier crossing replicate the experimental results to a reasonable degree of qualitative and quantitative fidelity.

## 2. RESULTS AND DISCUSSION

**2.1. Intramolecular KIEs.** The intramolecular KIE,  $k_{\text{HD}}/k_{\text{DH}}$  for the Chugaev elimination of  $D_1$ -**1** were measured by quantifying the resulting isotopic product distribution of  $D_1$ -**2** vs **2** (eq 1) using quantitative  $^1\text{H}$  NMR. The first letter in the subscript of  $k_{\text{HD}}$  indicates the transferred Hydron, and the second letter indicates the retained Hydron.

$$\text{KIE}_{\text{Intra}} = \frac{1^\circ\text{KIE}}{2^\circ\text{KIE}} = \frac{D_1\text{-}2}{2} \quad (1)$$

The Arrhenius-type plot of  $\ln(\text{KIE}_{\text{intra}})$  over  $1000/T$  is displayed in Figure 1. The temperature dependence of measured intramolecular KIE over 180 to 290 °C is  $k_{\text{HD}}/k_{\text{DH}} = 0.76 \exp(1.08 \text{ kcal}\cdot\text{mol}^{-1}/RT)$ . At 190 °C, the intramolecular KIE is  $2.45 \pm 0.04$ . Extrapolated to 25 °C, this KIE is 4.02. By



**Figure 1.** Arrhenius plot of  $\ln(\text{KIE}_{\text{intra}})$  vs  $1000/T$  for experimental and computational intramolecular KIE at 290, 270, 250, 230, 210, 200, 190, and 180 °C. The linear fit information is shown on the graph. The computational KIEs are computed by POLYRATE<sup>30</sup> using TST and CVT+SCT methods are also shown for comparison. These data are noteworthy for three reasons. First, nearly equivalent upward concave curvature is evident in the Arrhenius plots of both primary and secondary KIEs. As nonlinear, curved temperature dependence of  $\ln(\text{KIE})$  is one important indicator of tunneling,<sup>31</sup> our results indicate significant tunneling effect. Second, the magnitudes of the KIEs, especially secondary KIEs, extend beyond what might be expected in a reaction with behavior which is consonant with transition state theory. At the lower-temperature-end, both primary and secondary KIEs substantially exceed the values predicted by TST (solid lines in

definition, the intramolecular KIE for C–H cleavage from a methylene carbon is the ratio of the primary to the secondary KIE (eq 1). Even if one assumes a substantial  $2^\circ$  H/D KIE of approximately 1.2, which is typically associated with late, product-like transition states in processes like solvolysis, the primary H/D KIE at 25 °C could be assigned a maximal value of 4.83.<sup>26</sup>

The experimental intramolecular KIE results indicate: the ratio of the pre-exponential factors ( $A_{\text{HD}}/A_{\text{DH}} = 0.76$ ) is within the classical range ( $0.7 \sim \sqrt{2}$ ), the difference in the activation energy ( $\Delta\Delta E = \Delta E_{\text{DH}} - \Delta E_{\text{HD}} = 1.08 \text{ kcal}\cdot\text{mol}^{-1}$ ) is small, and the KIE values are moderate. According to classic works by Bell,<sup>7</sup> Stern,<sup>27</sup> and Weston,<sup>28</sup> these values do not make a convincing case that hydrogen tunneling is a dominant or even substantial factor in the Chugaev elimination.

It could be argued that the experimental data in Figure 1 possess a small degree of upward concave curvature. However, the confidence with which concavity may be imputed is questionable. It is noteworthy that the computed KIEs using TST method *exceed* experimental values by a significant margin, e.g., the computed intramolecular KIE at 190 °C is 2.82. This trend runs contrary to what is most frequently observed, where calculations based on TST with no tunneling correction typically yield smaller KIEs than experimental values, where tunneling is thought to have even the smallest influence.<sup>29</sup> Computational overestimation of intramolecular KIEs in the Chugaev elimination led us to interrogate this system further.

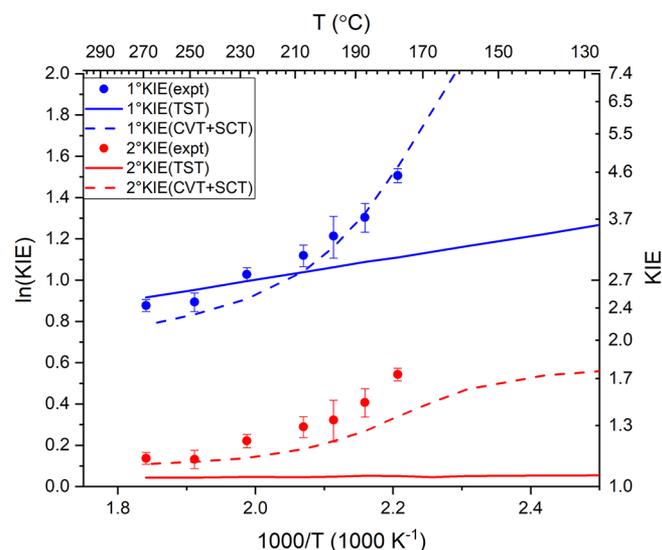
**2.2. Primary and Secondary KIEs.** For systems in which the rate- and product-determining steps are identical, the intramolecular and intermolecular KIEs can be used to solve for primary ( $1^\circ$ ) and secondary ( $2^\circ$ ) KIEs for hydrogen transfer reactions at a methylene group. The intermolecular KIE that reflects the relative elimination rates between **1** and  $D_1$ -**1** is a function of the  $1^\circ$  and  $2^\circ$  KIEs as shown in eq 2. Equations 1 and 2 can be solved simultaneously to yield expressions of the primary and secondary H/D KIEs (eqs 3 and 4), even though these values cannot be directly measured. (Refer to the Supporting Information, section III-(iv) for derivation.)

$$\text{KIE}_{\text{Inter}} = \frac{2}{1/1^\circ\text{KIE} + 1/2^\circ\text{KIE}} \quad (2)$$

$$1^\circ\text{KIE} = \frac{\text{KIE}_{\text{Inter}}(1 + \text{KIE}_{\text{Intra}})}{2} \quad (3)$$

$$2^\circ\text{KIE} = \frac{\text{KIE}_{\text{Inter}}(1 + 1/\text{KIE}_{\text{Intra}})}{2} \quad (4)$$

Extracting primary and secondary KIEs and using eqs 3 and 4 from 180 to 270 °C yields the Arrhenius plots shown in Figure 2. Values computed by POLYRATE<sup>30</sup> using TST and CVT+SCT methods are also shown for comparison. These data are noteworthy for three reasons. First, nearly equivalent upward concave curvature is evident in the Arrhenius plots of both primary and secondary KIEs. As nonlinear, curved temperature dependence of  $\ln(\text{KIE})$  is one important indicator of tunneling,<sup>31</sup> our results indicate significant tunneling effect. Second, the magnitudes of the KIEs, especially secondary KIEs, extend beyond what might be expected in a reaction with behavior which is consonant with transition state theory. At the lower-temperature-end, both primary and secondary KIEs substantially exceed the values predicted by TST (solid lines in



**Figure 2.** Arrhenius plots of  $\ln(\text{KIE})$  vs  $1000/T$  for experimental and computational primary and secondary KIEs. Experimental KIEs are plotted at 270, 250, 230, 210, 200, 190, and 180 °C. Computational results are calculated by POLYRATE with GAUSSRATE interface to Gaussian, with TST or CVT+SCT method, at the M06-2X/6-31+G(d,p) level of theory, with SMD heptane solvation.

(Figure 2). This result is due to the relatively greater impact tunneling has at lower temperatures. What is surprising is the magnitude of the secondary KIE at 180 °C (1.72)! To our knowledge, this is among the largest secondary KIE measured. The fact that this substantial secondary KIE occurs at far above room temperature and is exalted to a greater degree than the primary KIE at the same temperature challenges the notions that only transferred hydrogenic positions can dominate the tunneling event and that tunneling phenomena are influential only at room temperature or below.

Finally, POLYRATE Computation with canonical variational transition state theory (CVT) and small-curvature tunneling correction (SCT) reproduced the concaved temperature

dependence of the  $\ln(\text{KIE})$  vs  $1000/T$  (Figure 2), which will be discussed in the following sections.

**2.3. Analysis of the Rule of Geometric Mean.** The intermolecular KIE between **1** and  $D_2$ -**1** was measured to examine the rule of geometric mean (RGM). According to RGM, the KIE of a doubly labeled species should be the product of the primary and secondary KIEs (eq 5) when tunneling has little influence upon the reaction in question.<sup>32,33</sup> Substantial violation of RGM is suggestive of the presence of significant tunneling effect in the reaction.<sup>34,35</sup> At 190 °C,  $\text{KIE}_{\text{inter\_DD}}$  was measured to be  $6.6 \pm 0.1$ . At the same temperature,  $1^\circ \text{KIE} \times 2^\circ \text{KIE} = (3.67 \pm 0.25) \times (1.50 \pm 0.10) = 5.51 \pm 0.52$ . The product of primary and secondary KIEs is significantly different from the dideuterated KIE, deviating from RGM and thus consistent with tunneling.

$$\text{RGM: } \text{KIE}_{\text{inter\_DD}} = \frac{k_{\text{HH}}}{k_{\text{DD}}} = \frac{k_{\text{HH}}}{k_{\text{DH}}} \cdot \frac{k_{\text{HH}}}{k_{\text{HD}}} = 1^\circ \text{KIE} \times 2^\circ \text{KIE} \quad (5)$$

It is worth mentioning that **1** and its isotopologs are shelf stable at room temperature and undergo Chugaev elimination with reportedly no other side reactions when heated (see Supporting Information, section V). The mechanism of Chugaev elimination has been well established, in which the thion sulfur abstracts a  $\beta$ -hydrogen via a cyclic transition state in the rate-limiting step. Previous  $S^{32}/S^{34}$  and  $C^{12}/C^{13}$  KIE studies have ruled out the rearrangement pathway of the xanthate under the reaction conditions.<sup>36,37</sup> The  $^2\text{H}$  NMR on p S30 also indicate no  $\text{PhCH}_2\text{CH}_2\text{SCOSMe}$  isomer present. Nonetheless, based on the experiments conducted here, one cannot definitively eliminate the possibility that there could be alternative pathways available that lead to side reactions that were not detected.

**2.4. Computational Studies.** Initial efforts to model the temperature dependence of the  $1^\circ$  and  $2^\circ$  H/D KIEs associated with the Chugaev elimination failed to reproduce observed trends both qualitatively and quantitatively. If, as Figure 1 and Figure 2 suggested, the  $2^\circ$ KIE was exceptionally

**Table 1. Experimental and Computational Primary and Secondary KIEs for the Chugaev Elimination Reaction**

T (°C)	$1^\circ$ KIE					$1^\circ Q_t$		
	experimental	TST	CVT	TST+W <sup>a</sup>	CVT+SCT	$1^\circ Q_{t(\text{expt/TST})}^b$	$1^\circ Q_{t(w)}^c$	$1^\circ Q_{t(\text{SCT})}^d$
180	4.51 ± 0.15	3.03	2.25	3.30	4.70	1.49	1.085	2.02
190	3.67 ± 0.25	2.97	2.20	3.21	3.75	1.24	1.082	1.65
200	3.36 ± 0.34	2.89	2.15	3.12	3.19	1.16	1.079	1.42
210	3.06 ± 0.16	2.83	2.12	3.04	2.84	1.08	1.077	1.29
230	2.80 ± 0.09	2.71	2.04	2.90	2.48	1.03	1.072	1.16
250	2.44 ± 0.11	2.59	1.98	2.77	2.30	0.94	1.067	1.11
270	2.40 ± 0.07	2.50	1.92	2.65	2.19	0.96	1.063	1.09
T (°C)	$2^\circ$ KIE					$2^\circ Q_t$		
	experimental	TST	CVT	TST+W <sup>a</sup>	CVT+SCT	$2^\circ Q_{t(\text{expt/TST})}^b$	$2^\circ Q_{t(w)}^c$	$2^\circ Q_{t(\text{SCT})}^d$
180	1.72 ± 0.05	1.05	1.12	1.07	1.40	1.63	1.015	1.26
190	1.50 ± 0.10	1.05	1.11	1.06	1.31	1.42	1.014	1.18
200	1.38 ± 0.14	1.05	1.11	1.06	1.25	1.32	1.014	1.12
210	1.34 ± 0.07	1.05	1.11	1.06	1.20	1.28	1.013	1.08
230	1.25 ± 0.04	1.05	1.11	1.06	1.15	1.19	1.013	1.04
250	1.14 ± 0.05	1.04	1.11	1.06	1.13	1.09	1.012	1.02
270	1.15 ± 0.03	1.04	1.10	1.06	1.12	1.10	1.011	1.01

<sup>a</sup> $\text{KIE}^{\text{TST+W}}$  is KIE calculated with Wigner tunneling correction at TST level,  $k^{\text{TST+W}} = \kappa^{\text{W}} \cdot k^{\text{TST}}$ . <sup>b</sup> $Q_{t(\text{expt/TST})} = \text{KIE}^{\text{expt}} / \text{KIE}^{\text{TST}}$ . <sup>c</sup> $1^\circ Q_{t(w)} = \kappa_{\text{HH}}^{\text{W}} / \kappa_{\text{DH}}^{\text{W}}$ . <sup>d</sup> $1^\circ Q_{t(\text{SCT})} = \kappa_{\text{HH}}^{\text{SCT}} / \kappa_{\text{DH}}^{\text{SCT}}$ ;  $2^\circ Q_{t(w)} = \kappa_{\text{HH}}^{\text{W}} / \kappa_{\text{HD}}^{\text{W}}$ ;  $2^\circ Q_{t(\text{SCT})} = \kappa_{\text{HH}}^{\text{SCT}} / \kappa_{\text{HD}}^{\text{SCT}}$ .

exalted by tunneling, we would need to employ an approach that is inherently multidimensional. Toward that end, we employed POLYRATE<sup>30</sup> with GAUSSRATE<sup>38</sup> interface to Gaussian 09<sup>39</sup> using the M06-2X<sup>40</sup>/6-31+G(d,p)<sup>41</sup> functional to compute rate constants and transmission coefficients as a function of isotopolog. Results under the assumptions of transition state theory (TST) and canonical variational transition state theory (CVT)<sup>42,43</sup> were obtained. Multidimensional tunneling corrections were calculated using small-curvature tunneling (SCT)<sup>44,45</sup> option in POLYRATE. Some computational results are plotted in Figures 1 and 2, while more results are tabulated in Table 1. Three major observations can be drawn from the computational results.

First, CVT+SCT calculation reproduces experimental results in magnitude better than TST. As shown in Figure 2, compared to TST method, CVT+SCT method generates smaller 1° KIE values at higher temperatures, greater 1° KIE values at lower temperatures, and greater 2° KIEs, approaching more closely to the experimental values. Despite that SCT is the dominate factor, CVT without any tunneling correction starts to correct the rate constant toward the right direction (Table 1). CVT generates larger 2° KIE and smaller 1° KIE values than TST. The accurate approximation of the nontunneling rate constant is the first step toward the overall accuracy of the rate constant, as tunneling correction is a multiplicative term to the semiclassical rate constant calculated based on either TST or CVT. For example, rate constant calculated with CVT+SCT is the product of CVT rate constant ( $k^{CVT}$ ) and SCT transmission coefficient kappa ( $\kappa^{SCT}$ ) (eq 6). Canonical variational theory (CVT) is known to correct the conventional

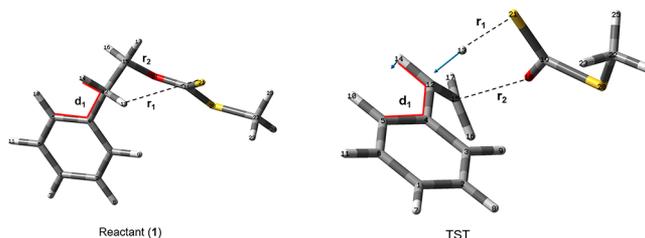
$$k^{CVT+SCT} = \kappa^{SCT} \cdot k^{CVT} \quad (6)$$

transition state theory (TST) by including zero-point energy and entropic effects. The location of the CVT transition state is the structure with the highest free energy while TST saddle point is the structure with the highest potential energy along the MEP.<sup>46</sup>

Second, CVT+SCT calculations successfully predict the concave-up trend of  $\ln(\text{KIE})$  vs  $1000/T$  for both primary and secondary KIEs. SCT tunneling correction reproduces the temperature dependence of primary and secondary KIEs better than other theories. As can be seen from Table 1, SCT correction term  $Q_{t(SCT)}$  is the major contributor to the inflation of the KIE values at lower temperatures, and is responsible to the curvature of the KIEs. As a comparison, the one-dimensional Wigner tunneling correction ( $Q_{t(W)}$ ) fails to predict either the magnitude or the trend for primary and secondary KIEs. The large curvature tunneling (LCT) calculation is useful for systems where the tunneling reaction path is close to a straight line from the reactant to the product valley.<sup>47,48</sup> However, the LCT method was not found to be the dominant tunneling mechanism in the Chugaev elimination reaction as it did not increase the transmission coefficient beyond SCT method. This is not unusual among other hydrogen transfer reactions.<sup>19,49,50</sup>

The third observation from POLYRATE calculation is that CVT+SCT calculation has limitations in predicting the tunneling effect on 2° KIEs. While experimental quantum amplification in 2° KIE is greater than that in 1° KIE ( $2^\circ Q_{t(expt/TST)} > 1^\circ Q_{t(expt/TST)}$ ), SCT calculation predicts the contrary ( $2^\circ Q_{t(SCT)} < 1^\circ Q_{t(SCT)}$ ).

To understand the tunneling behavior observed in the Chugaev elimination better, we examine the transition structures. The reactant 1 and the TST saddle point are optimized by Gaussian 09 at M06-2X/6-31+G(d,p) level of theory with SMD solvent model for heptane, and the structures are shown in Figure 3. The distance between the



**Figure 3.** M06-2X/6-31+G(d,p) structures of the reactant (1) and the TST saddle point with key metrics labeled.  $r_1$  and  $r_2$  denotes bond  $S_{21}-H_{13}$  and  $C_{15}-O_{18}$  respectively. Dihedral angle of  $H_{14}-C_{12}-C_4-C_5$  is outlined in red and labeled as  $d_1$ . The displacement vector of the imaginary frequency ( $952.65i \text{ cm}^{-1}$ ) of the saddle point is shown by blue arrows. The CVT transition structure is very close to the TST saddle point (see Table 2).

transferred hydrogen ( $H_{13}$ ) and the acceptor sulfur ( $S_{21}$ ) is denoted as  $r_1$ . The distance between  $C_{15}$  and  $O_{18}$  that undergoes bond breaking is denoted as  $r_2$ . The dihedral angle involving the hydrogen donor carbon ( $C_{12}$ ) that undergo  $sp^3$  to  $sp^2$  rehybridization are denoted as  $d_1$ .  $r_1$ ,  $r_2$ , and  $d_1$  are labeled and highlighted in Figure 3, and the values are tabulated in Table 2.

**Table 2. Selected Bond Lengths, Bond Orders and Dihedral Angle in the Reactant (1), TST Saddle Point, and the CVT Transition Structure**

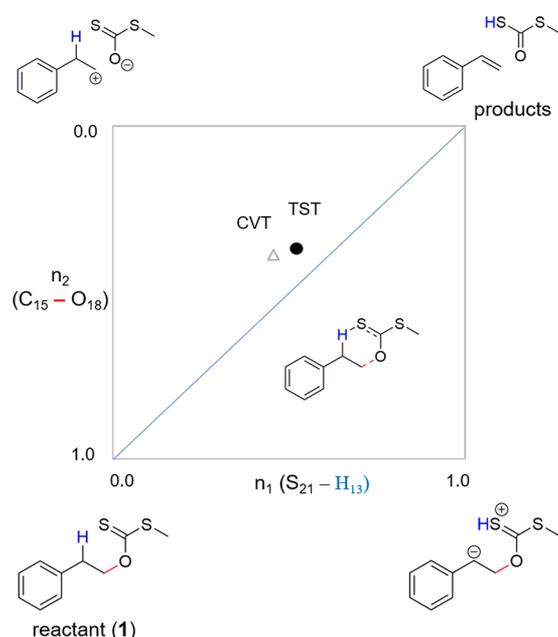
	$r_1$ (Å)	$r_2$ (Å)	$n_1^b$	$n_2^b$	$d_1$ (deg)
1	2.91	1.44	0.00 <sup>c</sup>	1.00	39.3
TST	1.74	2.04	0.52	0.37	3.4
CVT <sup>a</sup>	1.82	2.01	0.45	0.39	0.1

<sup>a</sup>CVT transition structure at 180 °C. <sup>b</sup>Bond order,  $n_i = \exp\left[\frac{(r_i^0 - r_i)}{C}\right]$  where  $r_i^0$  is the single bond length in the reactant or product in angstrom, and  $C = 0.6 \text{ Å}$ . <sup>c</sup>No bonding between  $S_{21}$  and  $H_{13}$  in 1.  $r_1^0 = 1.34 \text{ Å}$  in the product  $\text{CH}_3\text{SCOSH}$  is used in calculation of  $n_1$  for the saddle point and the CVT structure.

In addition, the key metrics of the CVT transition structure at 180 °C are also listed in Table 2. The bond order ( $n_i$ ) corresponds to a certain interatomic distance ( $r_i$ ) is calculated through Pauling relation<sup>51</sup> in eq 7, where  $r_i^0$  is the single bond length in the reactant or the product,  $r_i^0$  and  $r_i$  have a unit in angstrom, and  $C = 0.6 \text{ Å}$ .<sup>52</sup> The TST saddle point and CVT transition state are placed on a More O'Ferrall-Jencks diagram in Figure 4, in which the  $x$ -axis is the degree of bond formation of  $S_{21}-H_{13}$  bond, and the  $y$ -axis is the degree of  $C_{15}-O_{18}$  bond breaking.

$$n_i = \exp\left[\frac{r_i^0 - r_i}{C}\right] \quad (7)$$

The CVT transition structure is located more toward the reactant side compared to the saddle point on the reaction coordinate according to POLYRATE calculation, which is also



**Figure 4.** More O'Ferrall–Jencks diagrams for the Chugaev elimination of **1**. The solid circle represents the position of the TST saddle point, and the open triangle represents the position of CVT transition structure. Transition state positions are calculated from bond orders of  $S_{21}-H_{13}$  ( $n_1$ ) and  $C_{15}-O_{18}$  ( $n_2$ ) (values in Table 2). A diagonal line connecting the reactant and products is shown. The intermediate structures of two extreme scenarios: complete  $S_{21}-H_{13}$  breakage and complete  $C_{15}-O_{18}$  formation are shown in the lower-right and upper-left corners. The transition structure with partial  $S_{21}-H_{13}$  breakage and  $C_{15}-O_{18}$  formation is shown on the diagram.

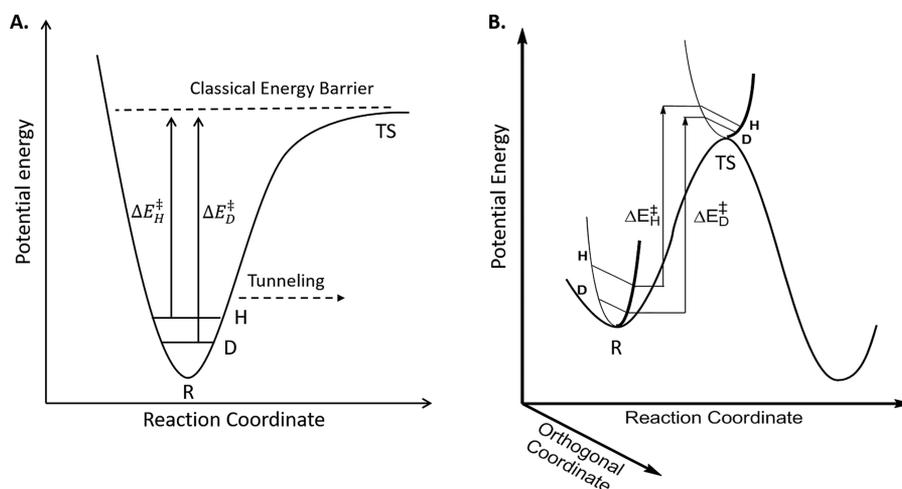
reflected in the More O'Ferrall–Jencks diagram (Figure 4). The CVT transition structure has less formation of  $S_{21}-H_{13}$  bond and less breakage of  $C_{15}-O_{18}$  bond (Figure 4 and Table 2), resembling the reactant more in terms of bond breaking/making. However, the dihedral angle  $d_1$  is more planar in the CVT transition structure (Table 2), which indicates a more

advanced  $sp^3$  to  $sp^2$  rehybridization at the  $C_{12}$  center, meaning the CVT structure resembles the product more than the TST structure in terms of rehybridization of  $C_{12}$ . These features align with the KIE values in Table 1: Compared to the TST method, the CVT method computes smaller  $1^\circ$  KIEs and greater  $2^\circ$  KIEs, corresponding to a less bond breaking/making but more rehybridized CVT transition structure.

The imaginary vibrational mode of the TST saddle point involves substantial movement for both the transferred hydrogen ( $H_{13}$ ) and the retained hydrogen ( $H_{14}$ ) while other atoms remain almost static (Figure 3 and animation file in the Supporting Information). For bound vibrational modes, the energy states are quantized (e.g., zero point energy) and proportional to the vibrational frequency (which is sensitive to both the “force constant” arising from the bonding environment and the mass of atoms involved in the mode); while for the unbound vibrational mode the deviation from classical behavior is manifested as the tunneling effect.<sup>7</sup> In KIE theories, quantum tunneling is predicted to have profound impact on  $1^\circ$  KIE but (often) negligible effect on  $2^\circ$  KIE, assuming the transferred hydrogen ( $1^\circ$  hydrogen) is involved in the reaction coordinate (unbound degree of freedom), while the motion of other atoms ( $2^\circ$  hydrogens) are contained in the orthogonal coordinates (Figure 5).<sup>15,16</sup> In this case, however, not only the transferred hydrogen but also the retained hydrogen is involved in the reaction coordinate. As a result, both  $1^\circ$  KIE and  $2^\circ$  KIE are significantly impacted by tunneling effect in the Chugaev elimination reaction. This coupling of  $2^\circ$  hydrogens to the reaction coordinate is not unprecedented, and some examples are discussed in the next section. In this study, the SCT model provides best agreement with experimental KIEs as it deals with the “corner-cutting” effect in which the system tunnel under the MEP, and it is generalized to deal with many-dimensional systems that have coupling of the reaction coordinate to orthogonal coordinates.<sup>45</sup>

## 2.5. Discussion in the Context of Related Research.

Measuring intramolecular KIE of the singly labeled substrate at a methylene group is a frequently applied method by organic



**Figure 5.** (A) Simple semiclassical model for primary/secondary H/D KIEs resulting from the difference in zero-point energies contained in the reaction coordinate. (B) H/D KIEs originating from zero-point energy differences at the reactant (R) and transition state (TS) in vibrational modes orthogonal to the reaction coordinate leading to a smaller activation energy  $\Delta E_H^\ddagger$  relative to  $\Delta E_D^\ddagger$ . Further rate enhancement occurs through tunneling into the classically forbidden region under the barrier, and it is related in part involved in the curvature of the reaction coordinate in the region of the TS as well as the mass of the hydrogen (H tunneling more effectively than D).<sup>15,16</sup>

chemists to study reaction mechanisms.<sup>53–56</sup> Unfortunately, this study implies that caution should be used when using intramolecular KIE to interpret reaction mechanisms, since some important information such as tunneling effects may not be revealed by intramolecular KIE, especially for elimination reactions.

There have been research efforts to study the tunneling effect upon secondary KIE in elimination reactions before this work. Saunders<sup>57</sup> reported a study on the tunneling contribution to the secondary isotope effect in 1984, and his model calculation study on E2 elimination reaction indicated that the extent of secondary KIE reflects the extent of the contribution of the nontransferred  $\beta$ -hydrogen (the hydrogen on the carbon adjacent to the leaving group) to the motion along the reaction coordinate. What we observe in Chugaev elimination reaction (an intramolecular *syn*- $\beta$ -elimination) agrees with what Saunders has found in E2 elimination. Intramolecular *syn*- $\beta$ -elimination reaction is similar to E2 reaction in every way except that the base is on the same molecule as the leaving group and the transition state structure is cyclic. The nontransferred  $\beta$ -hydrogen in both E2 and *syn*- $\beta$ -elimination reaction undergoes  $sp^3$  to  $sp^2$  rehybridization, resulting in significant motion in the reaction coordinate. More recent computational study by Borden<sup>58</sup> predicts significant tunneling effect in the reductive elimination of methane from an organo-platinum complex, resulting in violation of the rule of geometric mean and a large secondary KIE in the methyl group on the complex due to the motion of the methyl hydrogens in the imaginary vibrational mode. While the previous two works are computational studies, this work provides solid experimental evidence for quantum inflation of the secondary KIE.

Over 3 decades ago, a series of papers by Harold Kwart addressed a wide variety of intramolecular *syn*- $\beta$ -eliminations using the temperature dependence of intramolecular H/D KIEs.<sup>55,56,59–64</sup> Kwart (and more recently others<sup>65</sup>) have generated linear Arrhenius plots for all systems studied. One aspect of his work, specifically highlighted by Kwart, was the low magnitude of many of the intramolecular KIEs. This feature was attributed to a “bent” hydrogen transfer transition state. While such a scenario might explain Kwart’s results in part, work on other systems begin to illustrate that other factors may contribute to reduce intramolecular KIEs in *syn*- $\beta$ -eliminations. In particular, recent work in Meyer’s laboratory yielded an upward concave Arrhenius plot for the pericyclic *syn*- $\beta$ -elimination in the Swern oxidation of benzyl alcohol.<sup>66</sup> The intramolecular H/D KIE for the Swern oxidation of benzyl alcohol is abnormally low (2.53 at  $-78$  °C). The exceptionally low intramolecular H/D KIE for the Swern oxidation could be due to an abnormally low primary KIE, or an abnormally high secondary KIE, or both in concert. Unfortunately, in the Swern oxidation, the rate-determining step occurs prior to proton abstraction at the methylene in benzyl alcohol, so it is impossible to extract the information necessary to isolate the primary and secondary H/D KIEs. The Chugaev elimination, by contrast, has been shown to occur via rate-limiting hydrogen transfer.<sup>36,37</sup> Thus, we are able to extract primary and secondary KIEs from experimental intra- and intermolecular KIEs. Our results reveal that the quantum amplification of the secondary KIE is the cause of the suppression of the intramolecular KIE in Chugaev elimination. This finding suggests a possible explanation for Kwart’s results.

It should be mentioned that an E1 or E1cb mechanism is highly unlikely to give rise to the results presented here. While solvolysis mechanisms can give rise to large secondary H/D KIEs, the rate-limiting step for an E1 mechanism is formation of a carbocation at the carbon which is at the  $\beta$ -position relative to the benzene ring. The substantial  $1^\circ$  and  $2^\circ$  KIEs that manifest at the  $\alpha$ -position would be unlikely. Further, the rate-limiting step for an E1cb mechanism is proton abstraction. Aside from being chemically improbable, this mechanism would most plausibly be expected to yield a significant  $1^\circ$  KIE in conjunction with a negligible  $2^\circ$  KIE. These expectations along with previous  $^{13}\text{C}$  KIE studies<sup>36,37</sup> reinforce our assertions that neither of the above stepwise mechanisms are relevant to the Chugaev elimination studied here.

The contamination of intramolecular KIE by secondary KIE is less problematic for studies of enzymatic hydrogen transfer reactions, as the primary and secondary KIE are separable and can be measured directly in enzymatic reactions. This is because the substrates are chiral, and the enzymes are naturally stereoselective, e.g., alcohol dehydrogenase reaction,<sup>22</sup> reductive half-reaction of morphinone reductase involving NAD<sup>+</sup> and FMN,<sup>67</sup> etc. Experimental and computational evidence of significant tunneling in secondary KIE in enzymatic systems has been reported in NAD<sup>+</sup>/NAD-dependent dehydrogenases.<sup>20</sup>

### 3. CONCLUSION

In summary, this work presents insightful experimental and computational results that can serve as benchmark data for the development of more comprehensive hydrogen transfer theories. We have demonstrated three features of isotope effects in the Chugaev elimination. First is the seemingly classical Arrhenius plot corresponding to the intramolecular H/D KIE which masks pronounced quantum behavior. Second is the exalted secondary H/D KIE at 190 and 180 °C. Third is the observation of quantum behavior at high temperatures, some of which has been hinted at in the [1,5]-hydride shift and may be operative in other *syn*- $\beta$ -eliminations as a result of the cyclic nature of their transition structures and their substantial reaction barriers.<sup>68–70</sup>

The unimolecular, cyclic transition states which govern *syn*- $\beta$ -eliminations are likely to find parallels in hydrogen transfer steps in modern reactions of synthetic utility, e.g. those which implement (or suffer from)  $\beta$ -hydride eliminations and numerous intramolecular C–H functionalizations.<sup>53,71–75</sup> We hope that studies such as the one described here can both stimulate the development of new theoretical treatments of hydrogen transfer which can treat both primary and secondary motion in a quantum framework and help to encourage the search for such anomalous behaviors in synthetically and biologically relevant hydrogen transfer reactions.

### ■ ASSOCIATED CONTENT

#### 📄 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpca.9b00172.

Synthesis and characterization of **1**,  $D_1$ -**1**, and  $D_2$ -**1** experimental procedures for measuring KIEs, table of intra- and intermolecular KIE values, raw NMR integration data, energies and Cartesian coordinates of **1** and the transition structure, POLYRATE.dat file, and

POLYRATE output rate constants and transmission coefficients (PDF)

Animation of the transition vector (AVI)

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X.L. performed the KIE measurements, labeled substrate synthesis, and POLYRATE calculations. D.M.Y. provided constructive discussion and guidance in the major revision of the manuscript. M.P.M. designed the research project and performed Gaussian calculations.

### Notes

The authors declare no competing financial interest.

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